Optimization Of Mechanical Strength Of Rotary-Vane Engine

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Abstract. The article discusses a rotory-vane heat engine with a lever-cam mechanism motion conversion (an engine may be an internal combustion or external combustion). The output shaft of the engine adds drive torque from the working fluid pressure forces acting on the blade and the inertial moment of the forces of inertia of engine components. The mechanical strength of the motor is dependent on the magnitude and phase of these two torque. The purpose of the article is to determine the conditions under which mechanical strength is minimized.

Keywords: rotary-vane engine, lever-cam motion converter, external supply of heat.

I. INTRODUCTION

It is known that internal combustion engines have almost reached the limit of their energy and environmental characteristics. Therefore, it is necessary to develop engines of a new generation, constructed using other design schemes and implementing new thermodynamic cycles. To solve this problem rotor-vane machines with external heat supply can be used.

The Fig. 1 shows the engine component diagram.

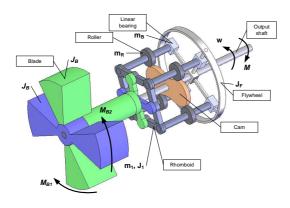


Fig. 1. Engine component diagram.

The engine consists of a driving part, motion converter, output shaft to the flywheel. The driving part consists of two coaxial rotors. Each rotor has the blade of the two diametrically opposite vanes (pistons). Also, the rotors have forks with two diametrically arranged fingers. The fingers are pivotally connected to the midpoints of the sides of the rhomboid links. The pressure of the working fluid between two adjacent vanes create moments on the blades $M_{B1} = -M_{B2}$, which depend on the output shaft angle.

The tops rhomboid are rollers. The rollers contact with a stationary profiled cam surface. Rhomboid is a motion transformer. It converts cyclic motion of the blade (within the angles ψ_{\min}, ψ_{\max} between them) into a unidirectional rotation of the output shaft. Thus transmitted torque from the blades. Torque transmission from rhomboid to the output shaft occurs via linear bearings.

To the output shaft of the engine is attached a load torque M.

The motion transformer uses a cam profile which in polar coordinates ρ , α is given by [1]:

$$\rho(\alpha) = 2l\sin(a+b\sin 2\alpha), \qquad (1)$$

where 2l – rhomboid link length, $a = \pi/4$; $b = \pi/4 - \psi_{min}/2$, where ψ_{min} – angle between the axes of the blades (usually $\psi_{min} = 50^{\circ} - 55^{\circ}$).

The vane angles change in a law:

$$\varphi_1(\alpha) = \frac{3\pi}{4} + \alpha + b\cos 2\alpha,$$

$$\varphi_2(\alpha) = \frac{\pi}{4} + \alpha - b\cos 2\alpha.$$
(2)

The dynamic system has one degree of freedom, because position of all elements is determined by one coordinate – output shaft angle α .

The motion equation of the system can be written in the form of Lagrange second kind equation:

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$$\frac{d}{dt}\frac{\partial T}{\partial \dot{\alpha}} - \frac{\partial T}{\partial \alpha} = M_{DT} - M, \qquad (3)$$

where T – the kinetic energy of the system; M_{DT} – the drive torque at the output shaft; M – the load torque.

Equation (3) can also be written as an kinetostatic equilibrium equation:

$$L_{Z} + M_{DT} - M = 0, (4)$$

where L_z – motor torque inertia forces applied to the output shaft:

$$L_{\rm z} = -\left(\frac{d}{dt}\frac{\partial T}{\partial \dot{\alpha}} - \frac{\partial T}{\partial \alpha}\right). \tag{5}$$

II. TORQUE OF INERTIA FORCES

The kinetic energy of the system includes the kinetic energy of the two blades T_B , four links of the rhomboid T_L , four rollers rhomboid T_R , four linear bearings with fingers T_{LB} , output shaft with a flywheel T_F :

$$T = T_B + T_L + T_R + T_{LB} + T_F.$$
 (6)

Rotors carrying blades, make a rotary motion around a fixed axis with angular velocity $\dot{\phi}_1, \dot{\phi}_2$. Kinetic energy of rotors carrying blades $T_B = J_B \left(\dot{\phi}_1^2 + \dot{\phi}_2^2 \right) / 2$, where $J_B = J_{B1} = J_{B2}$ – moments of inertia of the blades B1 and B2; $\dot{\phi}_1 = \dot{\alpha} \left(1 - 2b \cdot \sin 2\alpha \right); \quad \dot{\phi}_2 = \dot{\alpha} \left(1 + 2b \cdot \sin 2\alpha \right).$

$$T_{B} = \frac{1}{2}\dot{\alpha}^{2} 2 J_{B} \left(1 + 2b^{2} - 2b^{2}\cos 4\alpha \right).$$
(7)

Links rhomboid make plane motion, so their kinetic energy is the sum of the kinetic energies of the translational and rotational motions: $T_L = \left[m_1 \left(2V_{C1}^2 + 2V_{C2}^2 \right) + J_1 \left(2\Omega_1^2 + 2\Omega_2^2 \right) \right] / 2$, where m_1, J_1 – mass and moment of inertia about the central axis of one link; V_{C1}, V_{C2} – linear velocities of the centers; Ω_1, Ω_2 – angular velocities of the centers links 1 and 2. Given that

$$V_{C1}^{2} + V_{C2}^{2} = \dot{\alpha}^{2} l^{2} \left(2 + 4b^{2} - 4b^{2} \cos 4\alpha \right),$$

$$\Omega_{1}^{2} + \Omega_{2}^{2} = \dot{\alpha}^{2} \left(2 + 4b^{2} - 4b^{2} \cos 4\alpha \right)$$

$$T_{L} = \frac{1}{2} \dot{\alpha}^{2} 4 \left[\left(m_{1} l^{2} + J_{1} \right) \left(1 + 2b^{2} - 2b^{2} \cos 4\alpha \right) \right].$$
(8)

Four roller rhomboid has kinetic energy

$$T_{R} = \frac{1}{2}\dot{\alpha}^{2}8m_{R}l^{2}\left(1+2b^{2}-2b^{2}\cos 4\alpha\right),$$
 (9)

where m_R – mass of one roller.

The kinetic energy of the four linear bearings include of the kinetic energy of rotational motion $T_{rot} = \frac{1}{2}\dot{\alpha}^2 4J_{LB}$ and kinetic energy of translational motion $T_{tr} = \frac{1}{2}2m_{LB}\left(V_A^2 + V_B^2\right)$, where V_A, V_B – absolute tops speed rhomboid A and B; m_{LB}, J_{LB} – mass and moment of inertia of linear bearings. We obtain

$$T_{LB} = \frac{1}{2}\dot{\alpha}^{2} \cdot$$
(10)
$$\left[4J_{LB} + 8m_{LB}l^{2} \left(1 + 2b^{2} - 2b^{2}\cos 4\alpha \right) \right].$$

Flywheel together with the output shaft has a kinetic energy:

$$T_F = \frac{1}{2}\dot{\alpha}^2 J_F, \qquad (11)$$

where J_F – moment of inertia of flywheel together with the output shaft.

Summing up the values of the kinetic energy (7) - (11), we obtain the kinetic energy of the engine:

$$T = \frac{1}{2}\dot{\alpha}^2 J(\alpha), \qquad (12)$$

where $J(\alpha)$ – moment of inertia of engine referred to the output shaft. Moment of inertia includes fixed and variable components

$$J(\alpha) = J_c - J_v \cos 4\alpha. \tag{13}$$

There are J_c – permanent component; J_v – amplitude of the harmonic component with a period of $\pi/2$:

$$J_{C} = J_{F} + 4J_{LB} + 2(1+2b^{2}) \cdot \left[J_{B} + 2(m_{1}l^{2} + J_{1}) + 4l^{2}(m_{R} + m_{LB})\right];$$

$$J_{V} = 4b^{2} \left[J_{B} + 2(m_{1}l^{2} + J_{1}) + 4l^{2}(m_{R} + m_{LB})\right].$$
(14)
(15)

Expression for the given moment inertia of the engine is:

$$L_{z} = -\dot{\alpha}^{2} 2J_{v} \sin 4\alpha - \ddot{\alpha} \left(J_{c} - J_{v} \cos 4\alpha\right). \quad (16)$$

III. DEFINITION OF DRIVE TORQUE

Torque of gas pressure forces on the blades $M_{B1} = -M_{B2}$ transmitted through the cam motion transformer to the output shaft create a drive torque M_{DT} .

Work of drive torque at the elementary move $\delta A(M_{DT}) = M_{DT} \delta \alpha$ must be equal to work of torques on the vanes $\delta A(M_{B1}, M_{B2}) = M_{B1} \delta \varphi_1 + M_{B2} \delta \varphi_2$.

The angles of the vanes φ_1 and φ_2 increment

 $\delta \varphi_1 = \frac{d\varphi_1}{d\alpha} \delta \alpha \quad \text{and} \quad \delta \varphi_2 = \frac{d\varphi_2}{d\alpha} \delta \alpha, \quad \text{or using (2),}$ $\delta \varphi_1 = (1 - 2b \sin 2\alpha) \delta \alpha, \quad \delta \varphi_2 = (1 + 2b \sin 2\alpha) \delta \alpha.$ Moments on the vanes make work

 $\delta A(M_{B1}, M_{B2}) =$

$$[M_{B1} + M_{B2} + 2b(M_{B1} - M_{B2})sin2\alpha]\delta\alpha.$$

Are equal, so the work of the driving torque and torque on the vanes (including $M_{B2} = -M_{B1}$) we obtain an expression for the type of drive torque

$$M_{DT}(\alpha) = 4b M_{B1}(\alpha) \sin 2\alpha.$$
(17)

It is periodic function of α with period $T = \pi / 2$, which can be expanded in a Fourier series, and represented in the form

$$M_{DT}(\alpha) = M_0 + A_1 \cos 4\alpha + A_2 \cos 8\alpha + A_3 \cos 12\alpha + \dots +$$
(18)

$$B_1 \sin 4\alpha + B_2 \sin 8\alpha + B_3 \sin 12\alpha + \dots$$

where

$$M_{0} = \frac{1}{T} \int_{0}^{T} M_{DT}(\alpha) d\alpha,$$
$$A_{k} = \frac{2}{T} \int_{0}^{T} M_{DT}(\alpha) \cos k 4\alpha d\alpha,$$
$$B_{k} = \frac{2}{T} \int_{0}^{T} M_{DT}(\alpha) \sin 4k\alpha d\alpha,$$
$$(k = 1, 2, ...)$$

 M_0 is the average value of $M_{DT}(\alpha)$ for the period, which is a useful driving torque, which is the product of the angular velocity $\omega = \dot{\alpha}$ equal to engine power $N = M_0 \omega$. The remaining terms, as a harmonic function of the shaft angle $\alpha = \omega t$, are classified as disturbing moments with frequencies $4\omega, 8\omega, 12\omega,...$ The presence of these harmonics increases the mechanical strength of the engine.

Now we shall find condition under which the mechanical strength of the engine is minimized.

IV. TORQUE ON THE OUTPUT SHAFT

In the steady state acceleration $\ddot{\alpha} = 0$, the angular velocity $\dot{\alpha} = \omega = const$, and the angle of rotation of the shaft $\alpha = \omega t$.

Inertia moment L_z , defined by (5) takes the form $L_z = -\omega^2 2J_v \sin 4\omega t.$

Torque at the output shaft M_s equals the sum of the drive torque M_{DT} defined by (18), and the inertial moment L_z :

$$M_s = M_{DT} + L_z =$$

$$M_0 + \left(B_1 - 2J_V\omega^2\right)\sin 4\omega t + \tilde{M}(t),$$
(19)

where indicated

$$\tilde{M}(t) = A_1 \cos 4\omega t + A_2 \cos 8\omega t + \dots +$$

$$B_1 \sin 4\omega t + B_2 \sin 8\omega t + \dots$$

When the condition

$$B_1 - 2J_V \omega^2 = 0 \tag{20}$$

is compensation by inertia forces the most significant component of the disturbing drive torque $B_1 \sin 4\omega t$ and, consequently, to minimize the mechanical strength of the engine.

It follows that the design of its engine parameters should be selected so that the nominal angular velocity ω_n meet the condition (20)

$$\omega_n = \omega^* = \sqrt{\frac{B_1}{2J_V}}.$$
 (21)

V.RESULTS OF NUMERICAL SIMULATION

Initial data for the simulation: torque on output shaft $M_{DT} = M_0 = 1000 N \cdot m$, speed to 1500 rpm.

Torques on the blades and the drive torque as functions of the angle α are shown in Fig. 3.

Expansion $M_{DT}(\alpha)$ into a Fourier series (18) gives the following coefficients (in N·m):

 $M_0 = 1075,$

$$A_1 = -217, A_2 = -468, A_3 = -266, A_4 = -72,$$

$$B_1 = 1690, B_2 = 478, B_3 = 26, B_4 = -96.$$

The most significant component of the disturbing harmonic drive torque is $B_1 \sin 4\omega t$, amplitude of which is 1.5 times higher than the constant M_0 . Compensation of this harmonic by inertia forces must take place according to (21) at the angular speed

$$\omega^* = \sqrt{B_1 / 2J_V} = \sqrt{1690 / 2^* 0.184} = 68 c^{-1}$$

or at a frequency rotation $n^* = 30\omega^* / \pi = 650 \, rpm$.

Inertia moment $L_z = -\omega^2 2J_v \sin 4\omega t$ has an amplitude $A_{Lz} = 2J_v \omega^2$, parabolically depending on the speed (Fig. 4). Fig. 5 shows the dependence of the maximum torque on the shaft $(M_s)_{max}$ of the speed. Up to the speed n^* a decrease of $(M_s)_{max} = M_0 + |\tilde{M}|$ take place due to the compensation of the inertial forces components of the drive torque $B_1 \sin 4\omega t$. At speeds $n > n^*$ the torque $(M_s)_{max}$ increases parabolically under the action of the inertial forces.

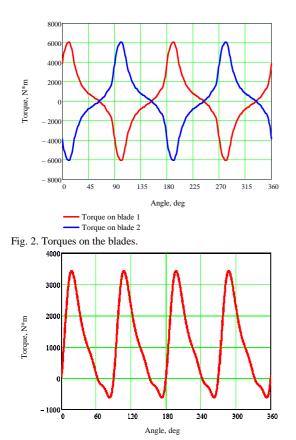
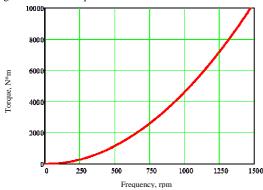


Fig. 3. The drive torque.



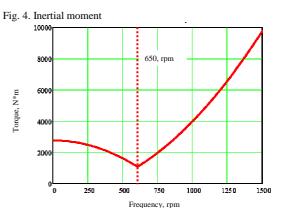


Fig. 5. Torque on the output shaft

VI. CONCLUSION

In this article, we have obtained the analytical expressions for motor driving and inertial moments of the single-module rotary-vane engine with lever-cam referenced to the output shaft, also the condition under which the forces of inertia completely compensate the most significant component of the disturbing drive torque and minimizes the mechanical strength of the engine.

The results of numerical simulation obtained on the example of the engine with internal combustion, which are close to the theoretical prediction.

This results can be used in the preliminary design of the rotary-vane engines.

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